



ELEN E3106/4106 Lecture 10

Metal-Semiconductor Junctions and Contacts

Outline

- Rectifying contacts
- Schottky barriers
- Schottky diodes
- Ohmic (non-rectifying) contacts

Assignments:

Reading: Streetman and Banerjee §5.7, C. Hu Ch. 4 Part III

Homework 4 due Fri. Oct. 10th by 5pm

Junction Recap

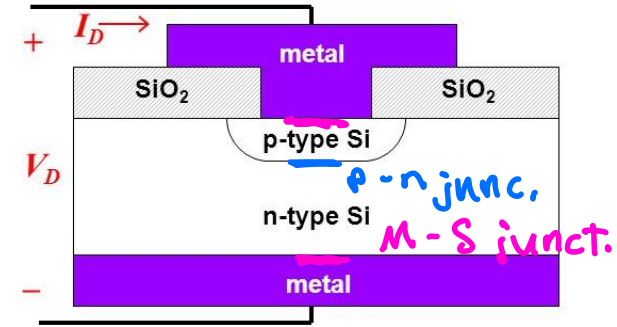
- Junction = boundary or interface between two types of materials.

Examples: $p-n$, p^+-n , n^+-p
 $n-n$, $p-p$

- Diode = semiconductor device whose principle of operation is based on the junction!
 - The physical structure has other practical components
 - Like contacts!
- Today we will be discussing another type of junction
 - Metal-semiconductor (M-S)

Physical structure:
(an example)

For simplicity, assume that the doping profile changes abruptly at the junction.



P-N junction



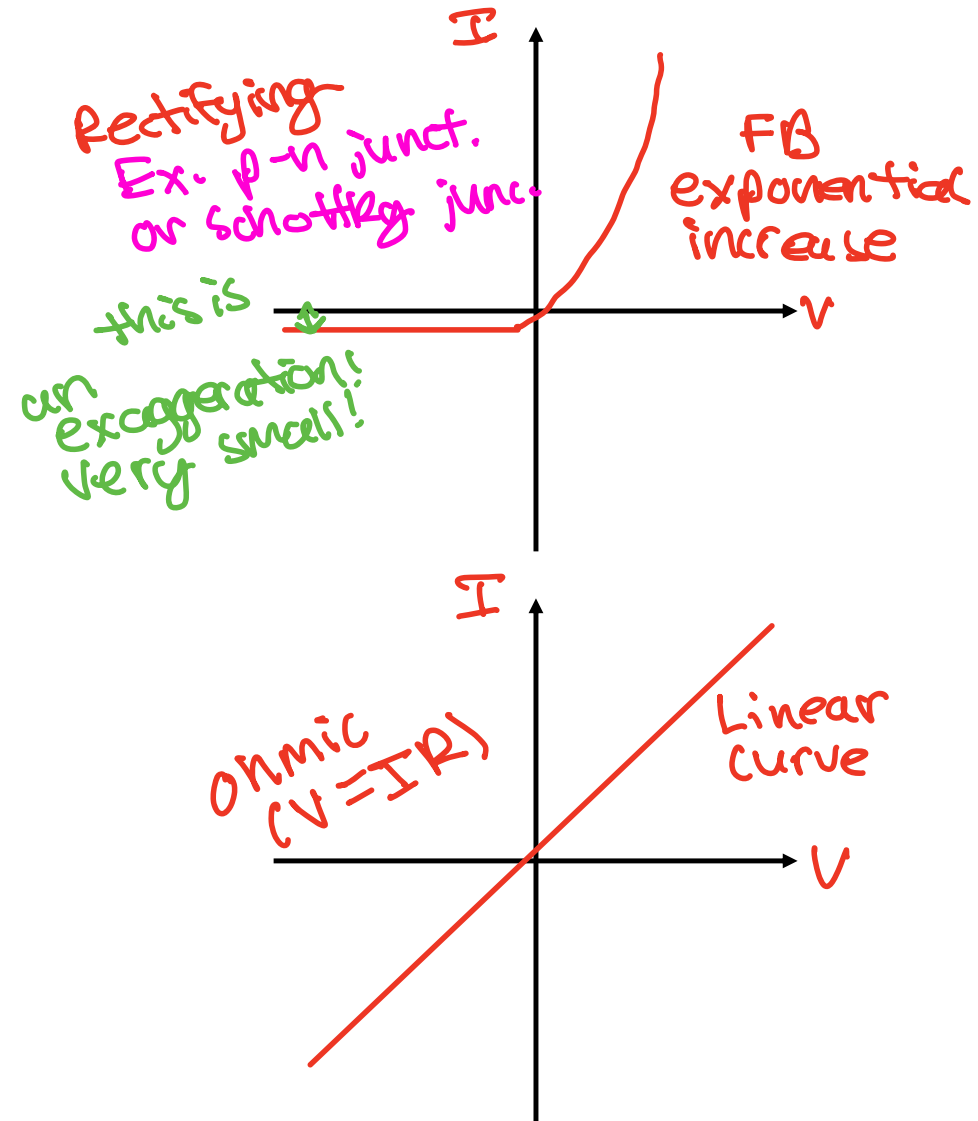
M-S junction



Any semiconductor

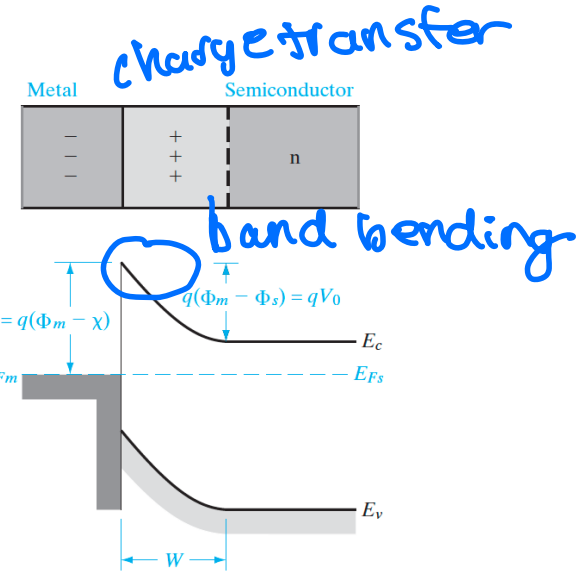
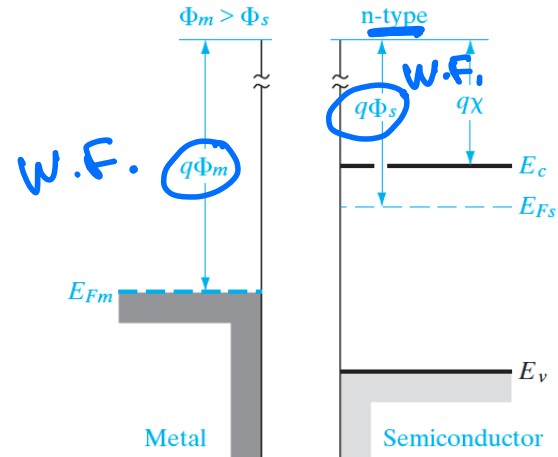
M-S Junctions

- Crucial to the operation of most semiconductor devices
- When metal and semiconductor are joined, two types of contacts can result:
 1. Rectifying (Schottky)
 2. Non-rectifying (ohmic)
- Rectifying means that current can only pass in one direction
- Non-rectifying means current can pass in both directions
- Question: Is a p-n diode rectifying?
yes!



Schottky Barrier Formation

- Recall: We discussed the metal work function $q\Phi_m$ (eV)
 - Energy needed to eject an e^- from E_F to vacuum outside the metal
- The semiconductor also has a work function $q\Phi_s = q\Phi_{vac} - E_F$
- Metal and semi brought in contact--> charge transfer occurs
- Because we are in equilibrium, E_F must align
- Let's consider when $\Phi_m > \Phi_s$ (n-type)
- Electrostatic potential of semiconductor must be raised (bands bend downward)

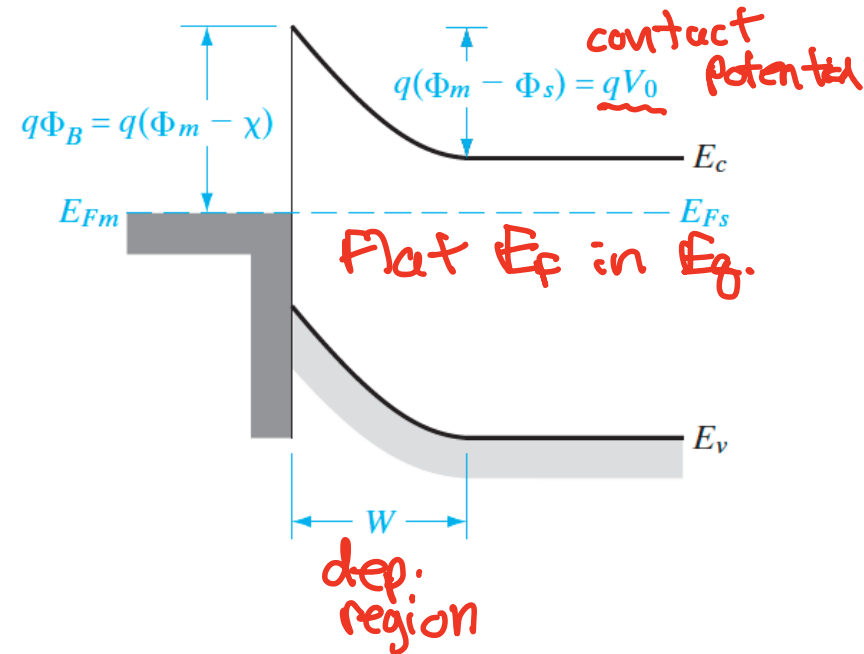
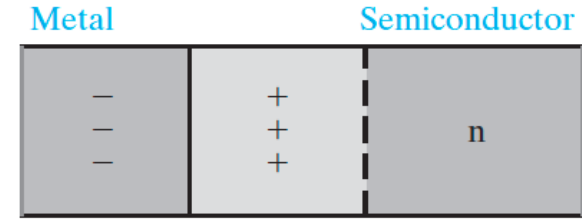


Schottky Barrier Depletion Region

- Depletion region w will form
 - (+) charge from uncompensated donors *(n-type)* N_d^+ in semiconductor
 - (-) charge is on the metal
- Built-in contact potential V_0 = $\Phi_m - \Phi_s$
- The case is similar to p^+-n approximation
 - (-) charge on metal is a very thin sheet to the left of junction

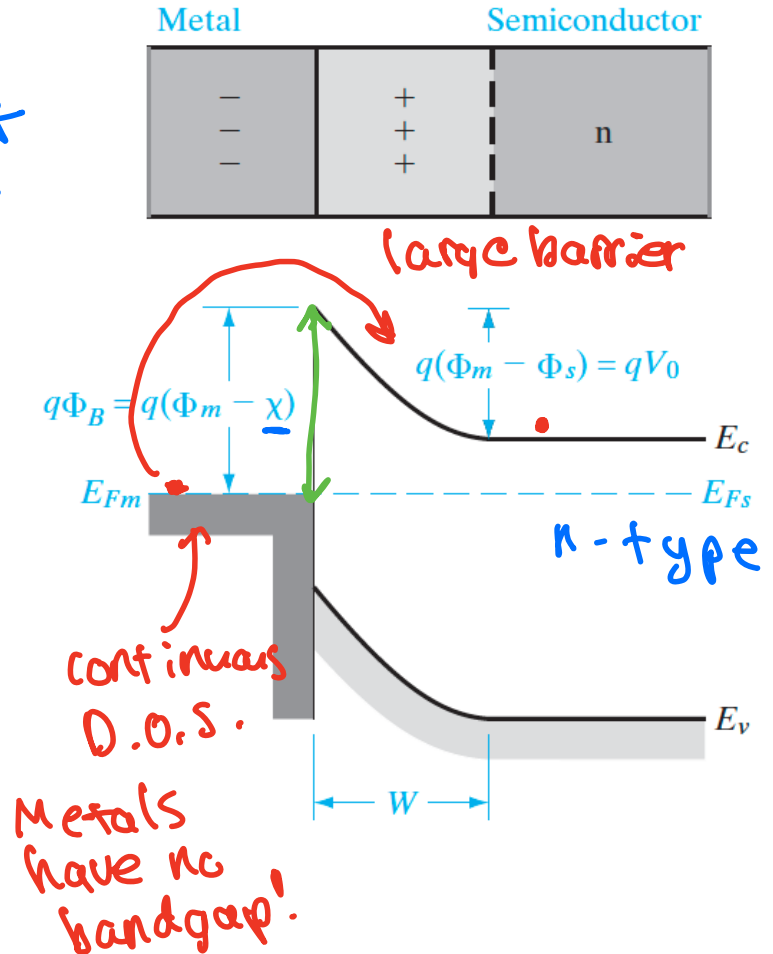
$$C_j = \left[\frac{2\epsilon V_0}{q} \left(\frac{1}{N_a} + \frac{1}{N_d} \right) \right]^{1/2}$$

$$C = A \frac{\epsilon_s}{W_{\text{dep}}}$$

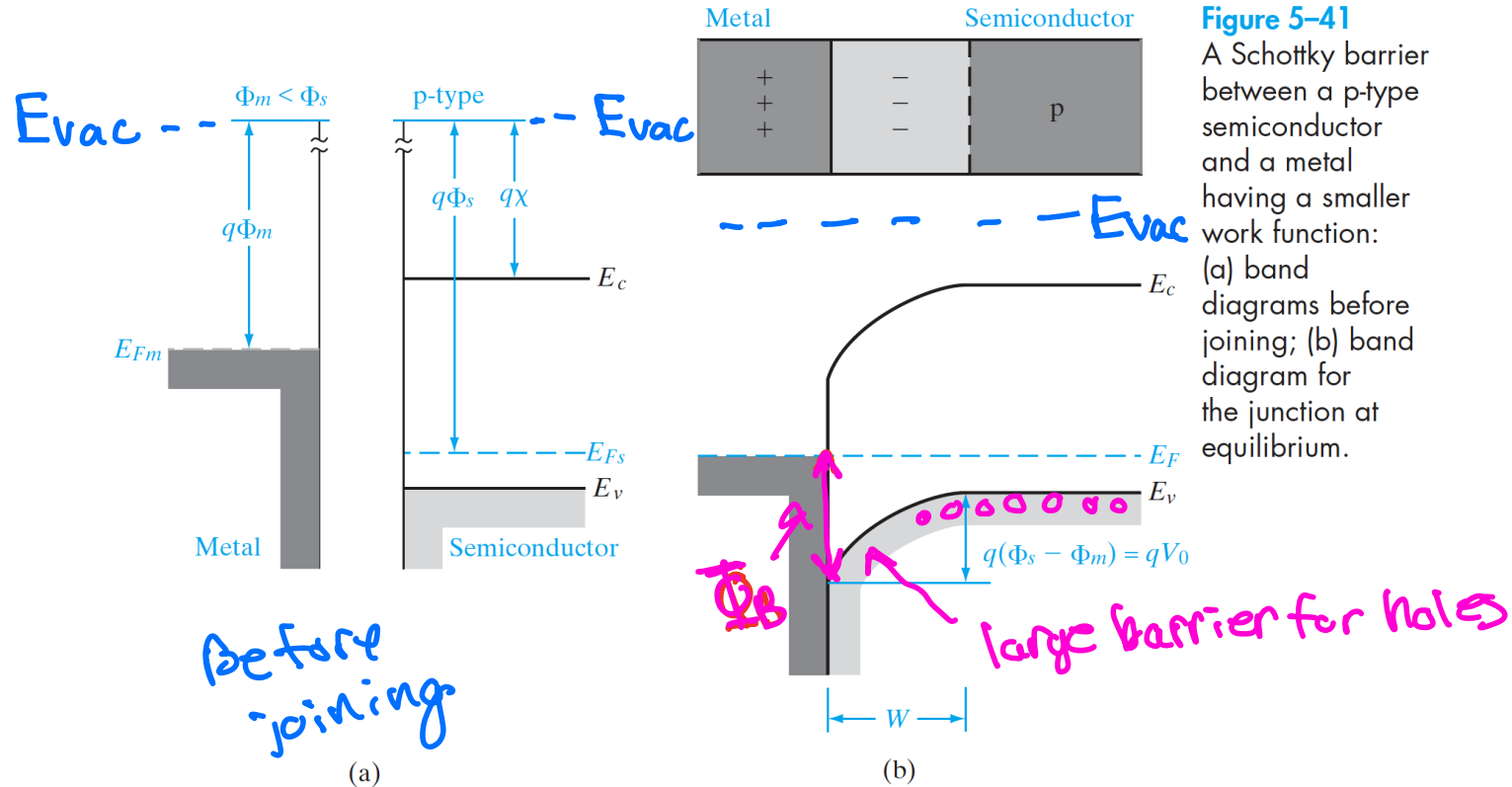


Schottky Barrier Height

- An energy barrier exists at the metal-semiconductor interface
- Schottky barrier height is denoted Φ_B (eV) *— barrier to current flow!*
- Φ_B depends on: (V)
 - Metal
 - Semiconductor
- The electron affinity, χ , is measured from the vacuum level to the semiconductor band edge $\chi = E_{vac} - E_c$
- $\Phi_B = \Phi_m - \chi$



Schottky Barrier on p-type semiconductor

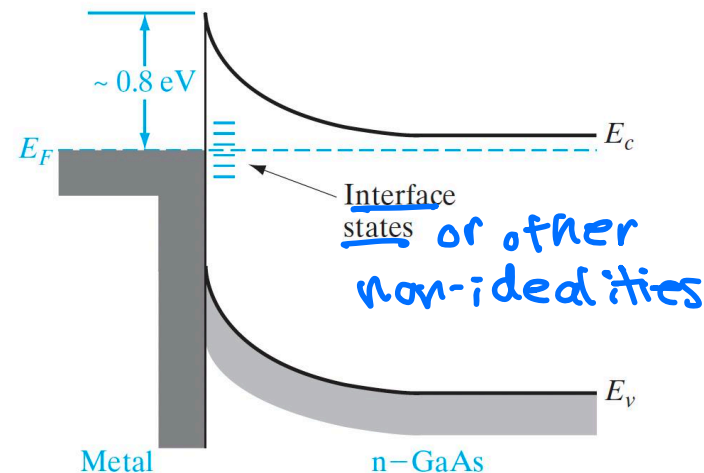


Fermi Level Pinning

- Why does Φ_B vary ^{metal} with the choice of the metal?
 $\Phi_B = \Phi_M - \chi$ ^{semiconductor}
- Clear trend that Φ_B increases with increasing metal work function ^(n-type)
- We would expect 1 eV change in Φ_m to result in 1eV change in Φ_B
 - Quantitatively not true!
- So far our discussion has looked at *ideal* Schottky barrier behavior
- In reality, the M-S junction is not ideal
 - Surface states due to incomplete covalent bonds
 - Interface roughness ^(measure w/AFM)
 - Thin interfacial layers
- These can pin the Fermi level at a certain position regardless of metal work function

TABLE 4-4 Measured Schottky barrier heights for electrons on N-type silicon (ϕ_{Bn}) and for holes on P-type silicon (ϕ_{Bp}). (From [7].)

Metal	Mg	Ti	Cr	W	Mo	Pd	Au	Pt
ϕ_{Bn} (V) ^{n-type}	0.4	0.5	0.61	0.67	0.68	0.77	0.8	0.9
ϕ_{Bp} (V) ^{p-type}		0.61	0.50		0.42		0.3	
Work Function ψ_M (V)	3.7	4.3	4.5	4.6	4.6	5.1	5.1	5.7



Using C-V Data to Determine Schottky Barrier Height

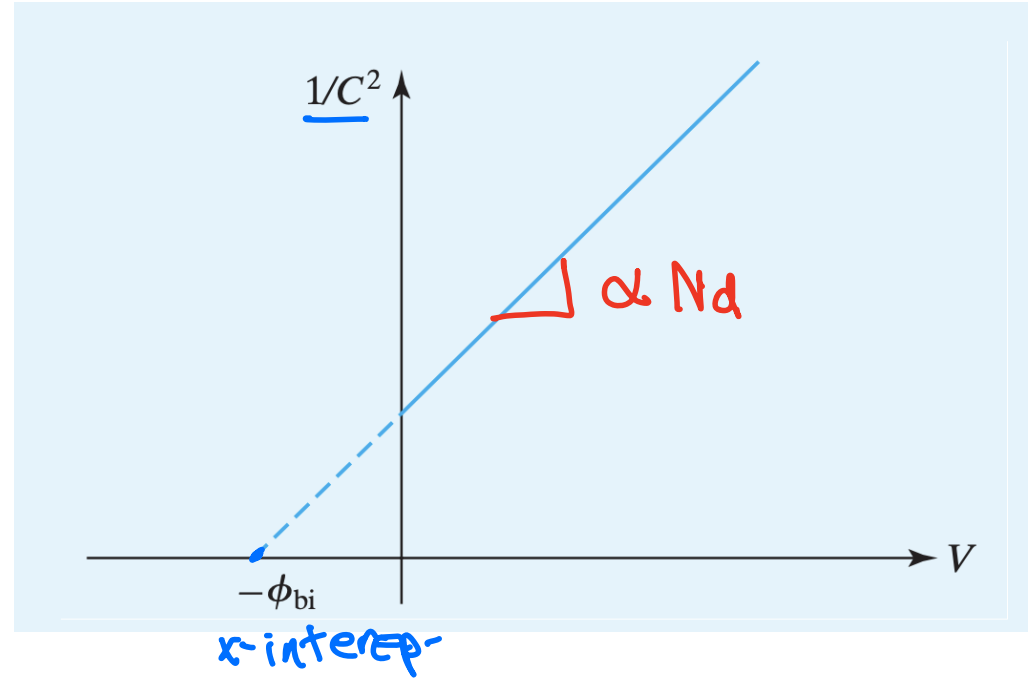
$$C = A \frac{\epsilon_s}{W_{\text{dep}}}$$

$$\frac{1}{C^2} = \frac{2(\phi_{bi} + V)}{q N_d \epsilon_s A^2}$$

$\Phi_{bi} = V_0$

$n\text{-semi}$

$$W_{\text{dep}} = \sqrt{\frac{2\epsilon_s(\phi_{bi} + V)}{q N_d}}$$



$$q\phi_{bi} = q\phi_{Bn} - (E_c - E_F) = q\phi_{Bn} - kT \ln \frac{N_c}{N_d}$$

← DOS

Similar to
 p^+n junction

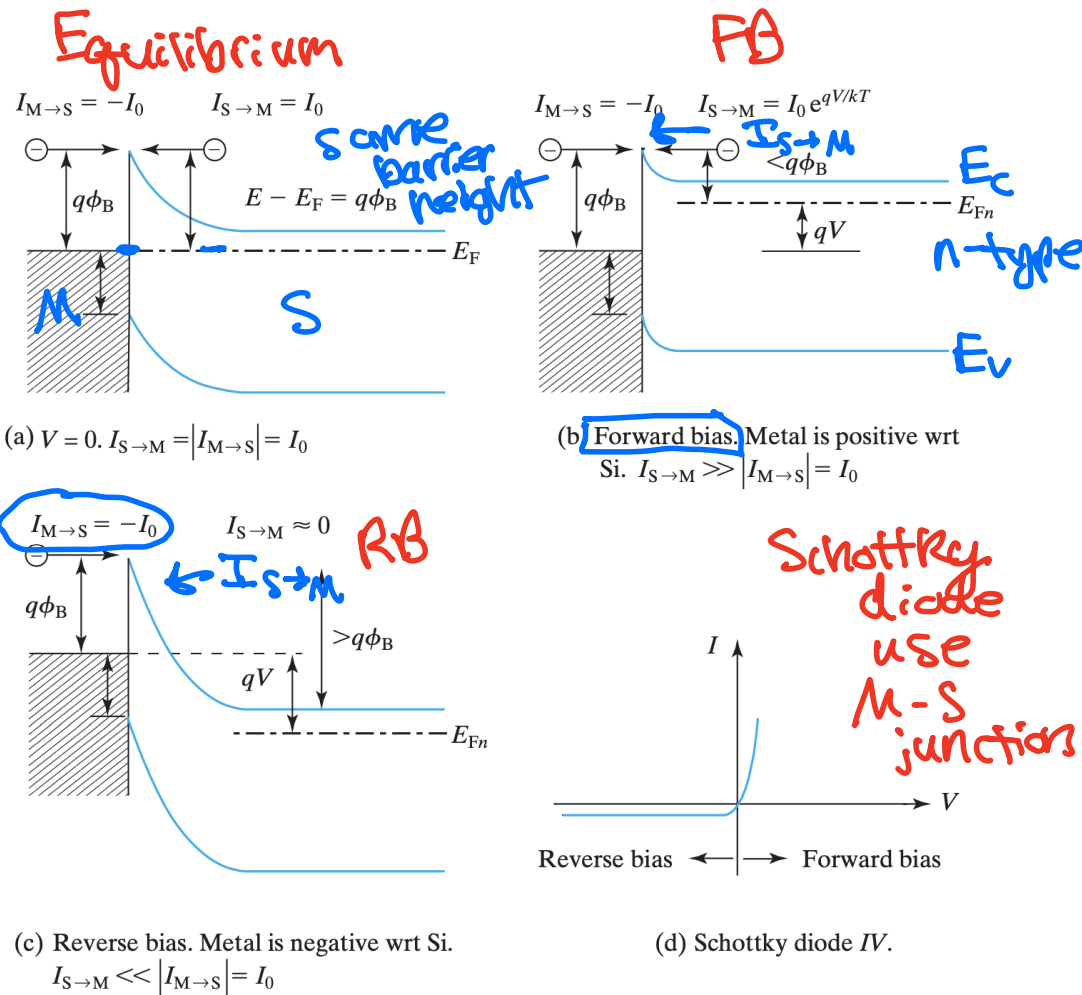
Explanation of Rectifying I-V Characteristics

- Rectifying, with easy current flow in the forward direction and little current in the reverse direction

$$I = I_0(e^{qV/kT} - 1)$$

- Forward current due to injection of majority carriers
- Left-traveling e^- have a thermal velocity
 - Thermionic emission theory

v_{th} (thermal velocity)
-x-dir.



Reverse Saturation Current

- Reverse saturation current is I_0
- Not as simply derived as in p-n diode, but what does it depend on?
Unaffected by bias (voltage)
Depends on Φ_B
- Is it affected by the bias voltage?

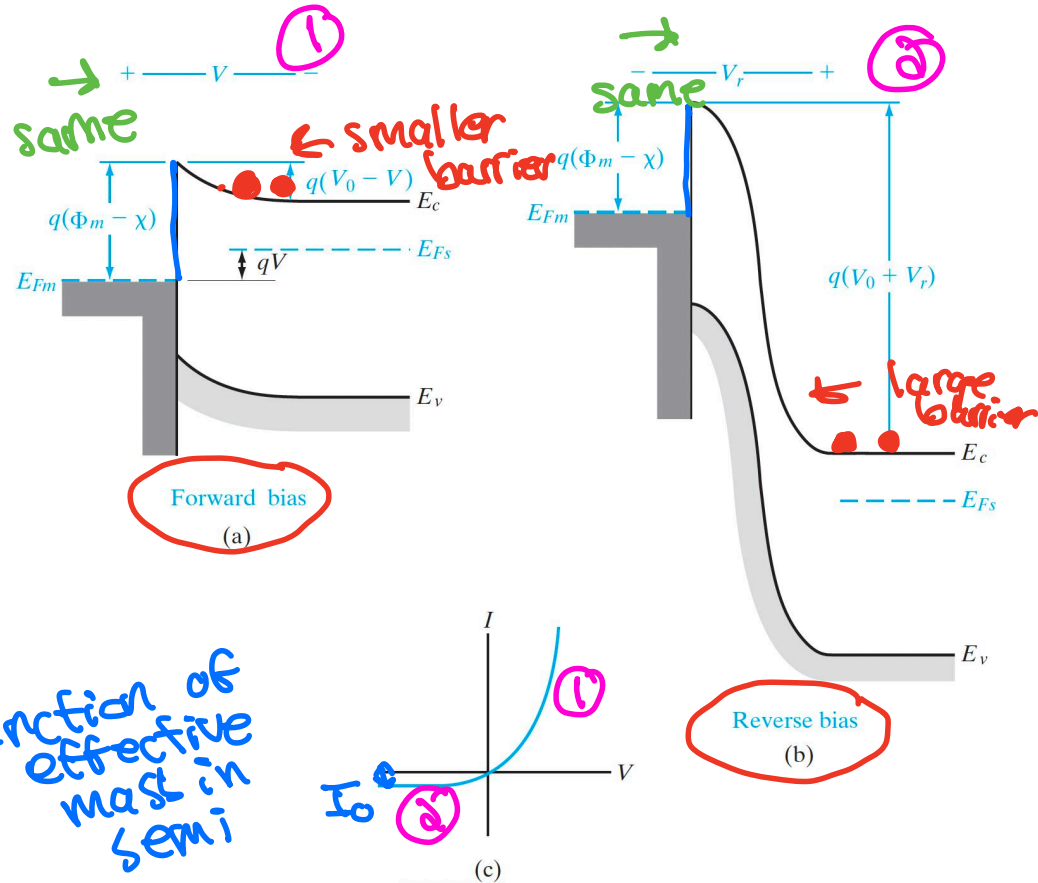
$$I_0 \propto e^{-q\Phi_B/kT}$$

$$I_0 = AKT^2 e^{-q\phi_B/kT}$$

- Richardson constant:

Units: $A/(cm^2 K^2)$ $K = \frac{4\pi q m_n k^2}{h^3}$

Fundamental parameter for thermionic emission



Schottky Diode Equations

- Desired: Φ barrier height, low I_0 (leakage current)
is high

- Ideal diode equation:

$$I = \underline{I_0}(e^{qV/kT} - 1)$$

- Taking into account non-idealities, forward current is:

$$I = A \overset{k}{\cancel{A}} T^2 e^{-q\Phi_B/kT} e^{qV/\overset{\text{ideality complex}}{\cancel{n}kT}}$$

Area

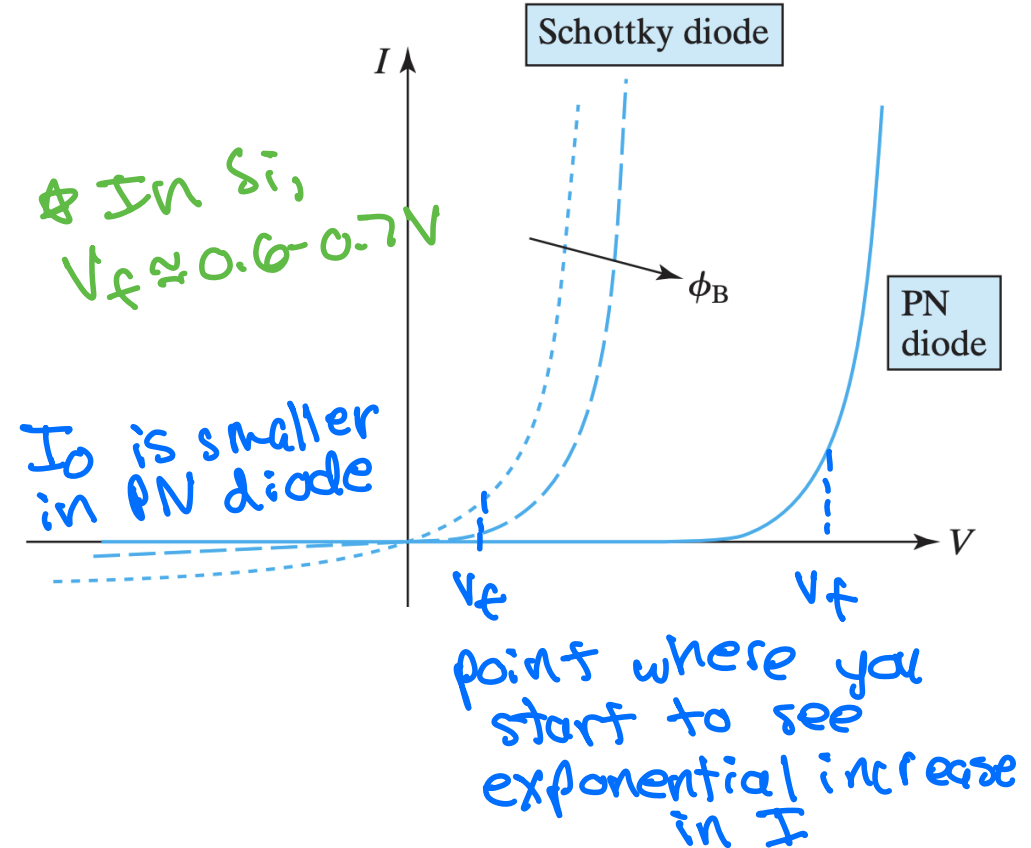
- Ideality factor n varies from 1-2 like p-n junction, but arises for different reasons

Schottky Diode Applications

- Advantages over p-n diode:
 - High-frequency properties and switching speed are generally better
 - Simpler fabrication process ✓
 - I_0 of a Schottky diode can be 10^3 – 10^8 times larger
 - Preferred in low-voltage and high-current applications
- Are there downsides to larger I_0 ?
Yes, larger power consumption in RB ($P = IV$)

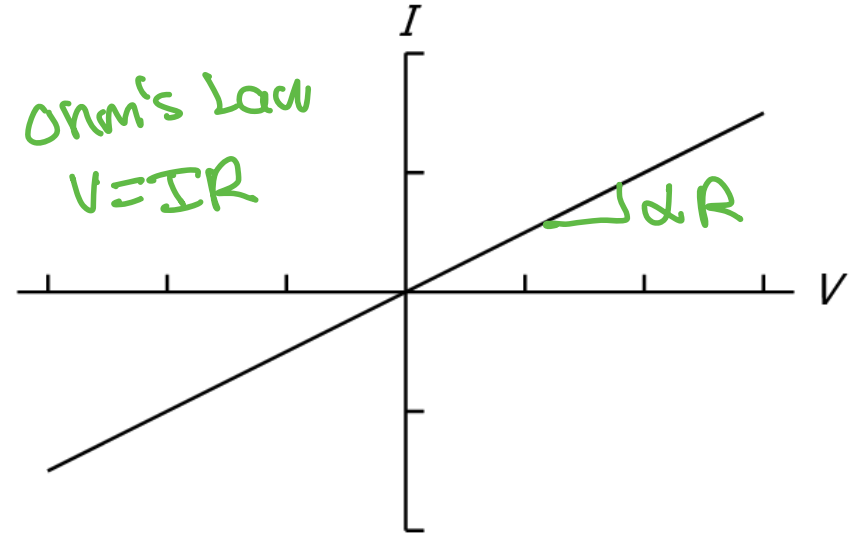
Share ideal diode eq.

$$I = I_0(e^{qV/kT} - 1)$$



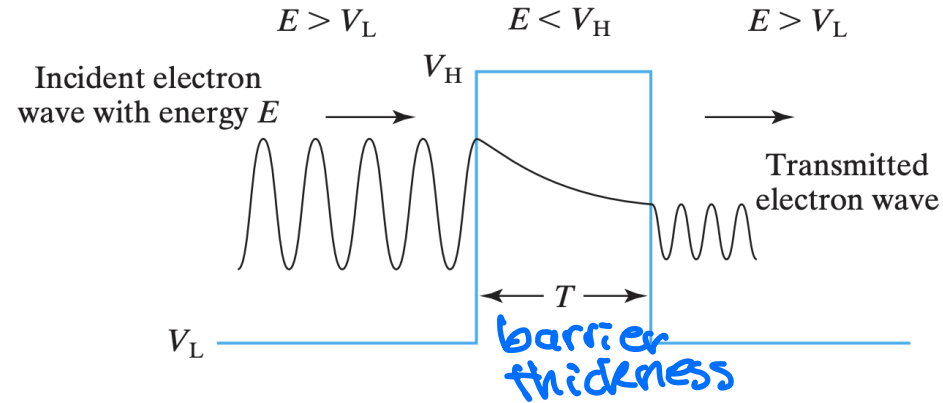
Non-Rectifying (Ohmic) Contacts

- M-S contacts are ohmic when we have linear I–V characteristic in both biasing directions
- What are the applications?
 - Contacting and interconnecting different regions of devices
- Desired properties of ohmic contacts:
 - Low resistance
 - No tendency to rectify
- 2 methods to formation:
- 1) Low barrier height
- 2) Heavily dope the semiconductor!
 - More practical!



Quantum Mechanical Tunneling

- There is a finite probability for e^- to tunnel through a potential barrier!
- Even if e^- doesn't have enough energy
- e^- with energy E arriving at barrier with high V_H
- Traveling wave decays
- Emerges from barrier with reduced amplitude
- Tunneling probability increases exponentially with decreasing barrier thickness



$$P \approx \exp\left(-2 \underbrace{T}_{\substack{\uparrow \\ \text{probability}}} \sqrt{\frac{8\pi^2 m}{h^2} (V_H - E)}\right)$$

Heavily Doped Ohmic Contacts

- With high dopant concentrations, barrier becomes thin

$$T \approx W_{\text{dep}}/2 = \sqrt{\epsilon_s \phi_{Bn} / (2qN_d)}$$

- High probability e- can pass through the barrier by tunneling
- Typical depletion layer width W of heavily doped Si?

Tens of angstroms

- Tunneling probability largely independent applied bias
 - Constant resistance (ohmic) contact

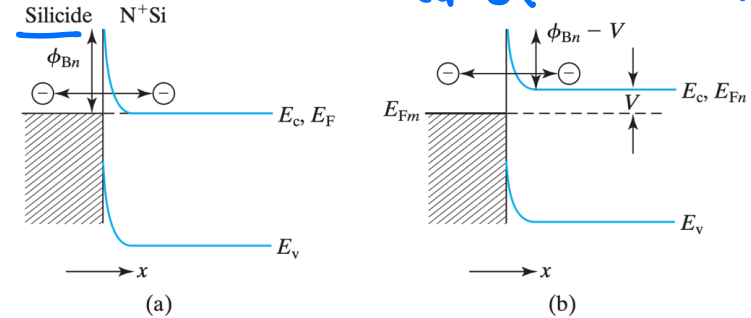
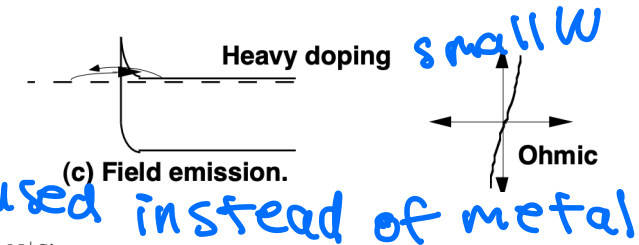
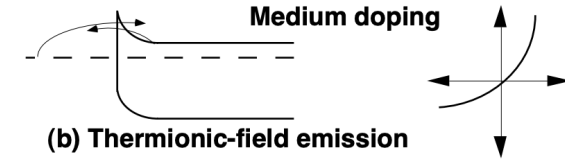
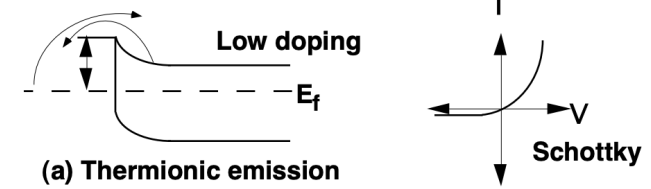


FIGURE 4-44 (a) Energy band diagram of metal- N^+ Si contact with no voltage applied and (b) the same contact with a voltage, V , applied to the contact.

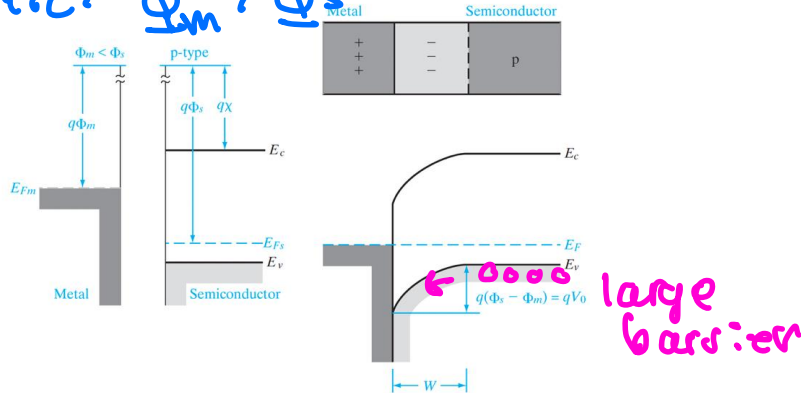
Summary of Schottky and ohmic contacts (low barrier height, method of formation #1):

- Metal/n-type semiconductor:

- Schottky: $\Phi_s < \Phi_M$
- Ohmic: $\Phi_M < \Phi_s$

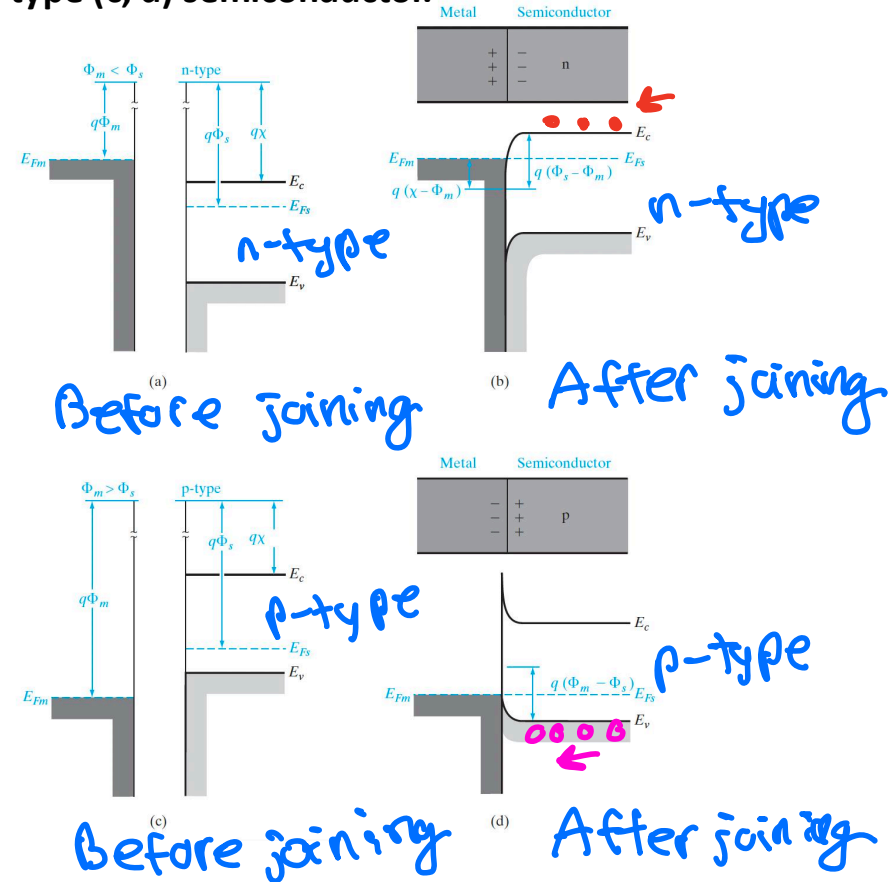
- Metal/p-type semiconductor:

- Schottky: $\Phi_s > \Phi_M$
- Ohmic: $\Phi_M > \Phi_s$



Schottky barrier on p-type semiconductor and a metal with a smaller work function.

Ohmic contacts on n-type (a, b) and p-type (c, d) semiconductor.



How do we find semiconductor work function?

- Φ_s : Work function of the semiconductor = $E_{vac} - E_F$
- χ : Electron affinity of the semiconductor = $E_{vac} - E_C$
- Looking at a band diagram, this means we can find Φ_s if we are given χ and the doping

- n-type semiconductor example:

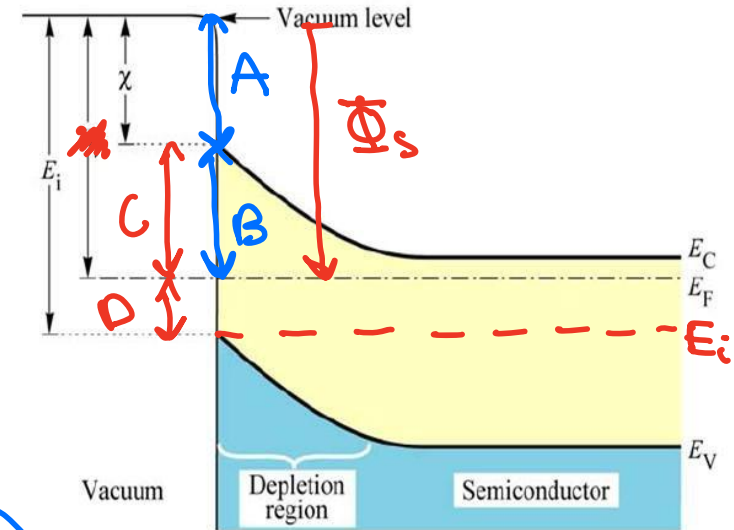
$$\begin{aligned}\Phi_s &= (E_{vac} - E_C) + (E_C - E_F) \\ \Phi_s &= (E_{vac} - E_C) + (E_C - E_i) - (E_F - E_i) \\ \Phi_s &= \chi + \frac{E_g}{2} - (E_F - E_i)\end{aligned}$$

- Recall: Equilibrium carrier concentrations (cm^{-3}):

$$n_0 = n_i e^{(E_F - E_i)/kT}, p_0 = n_i e^{(E_i - E_F)/kT}$$

Rearrange

$$\Phi_s = \chi + \frac{E_g}{2} - kT \ln \left(\frac{n_0}{n_i} \right) \quad (\text{n-type})$$



Contact Resistance

- Desired: low contact resistance
- An important property of contacts is the specific contact resistance, R_c
 • Units: $\Omega \cdot \text{mm}$ $R (\Omega)$
- A measure of how easily current flows through a contact
- For an n-type semiconductor:

$$R_c \equiv \frac{V}{J} = \frac{2 \cdot e^{H\phi_{Bn}/\sqrt{N_d}}}{qv_{thx} H \sqrt{N_d}}$$

A/mm
 $\propto e^{H\phi_{Bn}/\sqrt{N_d}}$

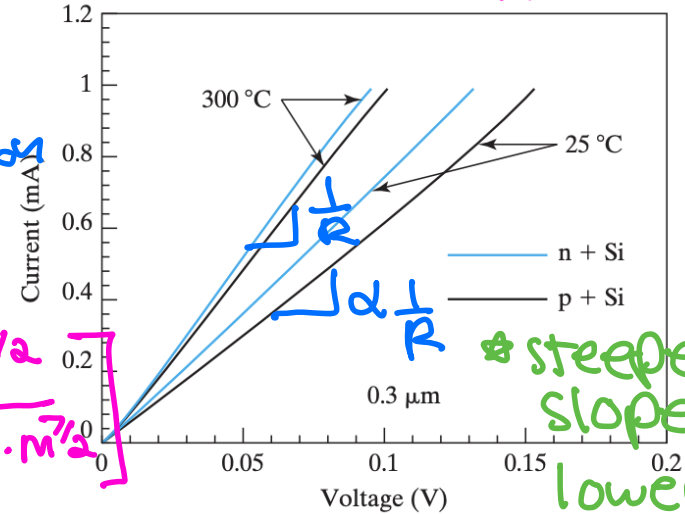
in x-direction
 $v_{thx} = -\sqrt{2kT/\pi m_n}$

$H \equiv \frac{4\pi}{h} \sqrt{(\epsilon_s m_n)/q}$

H is the

- Recall, v_{th} is the thermal velocity
- Same thing true for p-type, but we replace N_d, ϕ_{Bn}, m_n by N_a, ϕ_{Bp}, m_p

* H is not dimensionless but treated as material-dependent constant



I-V characteristics of a 0.3 micron diameter TiS_2 contact on n+ and p+ Si.

* steeper slope = lower resistance

tunneling factor, acts as proportionality term connecting barrier height and doping concentration